

Effective wavelength as a universal parameter of hyperspectral light radiance upwelling from the sea

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ABSTRACT

This paper presents experimentally obtained relationships between effective wavelength and a number of important seawater parameters such as dominant wavelength, different color indices, downward diffuse attenuation coefficient, absorption coefficient, seawater effective optical thickness, and a chlorophyll content. A number of regression relationships that connect hyperspectral measurements with listed bio-optical properties are proposed and discussed.

Keywords: marine optics, scattering, sea water, optical properties

1. INTRODUCTION

Processing of optical satellite data corrected from atmospheric effects usually consists in correlating different combinations of upwelling radiances with optical and biological parameters of seawater. When the number of channels is small it is possible to correlate almost all combinations of radiances with bio-optical properties. In processing hyperspectral data the number of channels is so large that it is impossible to effectively utilize all hyperspectral information using only radiances at certain wavelengths. In order to effectively process all hyperspectral information we need to introduce a new set of integral parameters that are related to all spectral channels. One of the simplest parameter is a normalized first moment of visible spectrum in relation to wavelength. We name this parameter as an effective wavelength of upwelling hyperspectral radiance.

This paper discusses different experimentally obtained relationships between effective wavelength and a number of important seawater parameters such as CIE dominant wavelength, different color indices, downward diffuse attenuation coefficient, absorption coefficient, seawater effective optical thickness, and a chlorophyll concentration. A number of regression relationships that connect hyperspectral measurements with listed bio-optical properties are proposed and discussed. It was shown that an effective wavelength is a universal parameter that connects hyperspectral optical measurements with a number of optical properties and effective chlorophyll concentration.

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1. INTRODUCTION

As an integral parameter of multispectral quantity $\Phi(\lambda)$, such as upwelling radiance $L(\lambda)$, irradiance $E(\lambda)$, remote sensing reflection coefficient $r(\lambda)$, or diffuse reflection coefficient $R(\lambda)$, etc. let us choose the following value,

$$\lambda_{eff}(\Phi) = \int_{\lambda_1}^{\lambda_2} \lambda \Phi(\lambda) d\lambda / \int_{\lambda_1}^{\lambda_2} \Phi(\lambda) d\lambda. \quad (1)$$

It is convenient to name this parameter as an effective wavelength over spectrum $\Phi(\lambda)$ in the range of $\{\lambda_1, \lambda_2\}$. If $\Phi(\lambda) = E_{\uparrow}(\lambda)$, λ_{eff} has a physical meaning: it is reversely proportional to the average energy of quanta in the range of $\lambda_1 \leq \lambda \leq \lambda_2$, i.e. $\lambda_{eff}(E_{\uparrow}) = hcQ/W$, where

$$Q = \frac{1}{hc} \int_{\lambda_1}^{\lambda_2} \lambda E_{\uparrow}(\lambda) d\lambda \quad (2)$$

is quantum irradiance from below, and

$$W = \int_{\lambda_1}^{\lambda_2} E_{\uparrow}(\lambda) d\lambda \quad (3)$$

is total energetic irradiance, h is the Plank's constant, c is the speed of light in vacuum. In the case of $\lambda_1 = 350$ nm, $\lambda_2 = 700$ nm, that we will use in this paper, and $\Phi(\lambda) = E_{\uparrow}(\lambda)$ Eq. (1) determines the effective wavelength of photo-synthetically active radiation. Because spectral shapes of $E(\lambda)$ and $L(\lambda)$ are identical [1], $\lambda_{eff}(E_{\uparrow}) \equiv \lambda_{eff}(L_{\uparrow})$. Numerical estimates show that normalizing $E(\lambda)$ and $L(\lambda)$ by downward irradiance $E_{\downarrow}(\lambda)$, which corresponds to $\Phi(\lambda) = R(\lambda)$ and $\Phi(\lambda) = r(\lambda)$, changes values of λ_{eff} less than 0.5%. Analysis of data collected during our joint shipborn and airborne experiments shows that effective wavelength weakly varies with the change in lightning conditions, and optical properties of atmosphere. At the same time λ_{eff} is strongly coupled with inherent optical properties of the water mass, light field parameters of the water, and concentrations of dissolved and suspended matter [2].

2. CONNECTION BETWEEN EFFECTIVE WAVELENGTH OF UPWELLING RADIATION AND LIGHT FIELD PARAMETERS OF SEAWATER

Being an integral property of the light upwelling from the seawater, specifically, being a wavelength averaged over light spectrum upwelling from the sea, the effective wavelength λ_{eff} is particularly sensitive to the highly energetic components of sea light that determine the color of the water mass. Comparison of computed values of effective λ_{eff} and dominant λ_D [3, 4] wavelengths revealed strong coupling between these properties. This coupling may be expressed by the following regressions:

$$\lambda_D = 213.6 + 0.56 \lambda_{eff}, \quad 459 < \lambda_{eff} < 497 \text{ nm}, \quad (N = 24, r = 0.96), \quad (4)$$

$$\lambda_D = -1038.4 + 3.07 \lambda_{eff}, \quad 498 < \lambda_{eff} < 521 \text{ nm}, \quad (N = 25, r = 0.93). \quad (5)$$

The color index, which was widely used in hydro-optical measurements since its introduction by Jerlov [5], also correlates nicely with the effective wavelength:

$$\lg J(\lambda_n, \lambda_k) = -A_1 + B_1 \lambda_{eff}, \quad (6)$$

$$\lg \lambda_{eff} = A_2 + B_2 J(\lambda_n, \lambda_k), \quad (7)$$

here and everywhere in this paper \lg means a logarithm of base 10.

Equations (4) and (5) are given for different combinations of wavelength pairs $\{\lambda_n, \lambda_k\}$, and coefficients A_1 , B_1 , A_2 , and B_2 in Eqs. (4)-(5) are given in Tables 1a, 1b, 2a, and 2b, respectively.

Not all color indices $J(\lambda_n, \lambda_k) = L_{\uparrow}(\lambda_n) / L_{\uparrow}(\lambda_k)$ correlate with λ_{eff} equally well, but in all cases presented in Tables 1 and 2 the correlation coefficients are larger than 0.8. For majority of the pairs presented here the correlation coefficient exceeds 0.9 and reaches maximum of $r = 0.996$ for $J(580, 420)$ [6].

In spite of the fact that the choice of wavelengths λ_n and λ_k is usually related to the maximums of phytoplankton absorption, there is no general agreement concerning the values of these wavelengths. This causes difficulties in comparing results of color index measurements obtained by different authors. The joint use of Eqs. (6) and (7) allow us to compare and generalize such previously incomparable measurements.

Some optical classifications of water masses [7-10] use spectral distribution of downward diffuse attenuation coefficient $k_{\downarrow}(\lambda)$ as their basis. This is possible because the $k_{\downarrow}(\lambda)$ averaged over thick subsurface layer of water

Table 1a. Numerical values of coefficient A_1 in Eq. (6).

$\lambda_n \setminus \lambda_k$	410	420	430	440	450	460	470
470	3.25	2.94	2.65	2.15	1.18	0.5	-
480	4.06	3.75	3.47	2.95	1.97	1.39	0.78
490	4.66	4.36	4.05	3.56	2.57	1.98	1.4
500	6.12	5.81	5.51	5.00	4.02	3.43	2.84
510	7.68	7.4	7.09	6.59	5.6	5.02	4.45
520	9.65	9.29	8.98	8.49	7.52	6.95	6.34
530	10.30	10.18	9.71	9.41	8.38	7.76	7.21
540	11.17	10.73	10.55	10.05	9.95	8.49	7.89
550	12.25	11.89	11.63	11.20	10.18	9.81	9.05
560	13.32	12.98	12.71	12.15	11.21	10.21	10.02
570	13.43	13.19	12.94	12.45	11.47	10.87	10.32
580	14.13	13.88	13.61	13.09	12.12	11.50	10.92
590	13.35	13.01	12.69	12.21	11.31	10.60	9.59

Table 1b. Numerical values of coefficient $B_1 \cdot 10^3$ in Eq. (6).

$\lambda_n \setminus \lambda_k$	410	420	430	440	450	460	470
470	6.7	6.1	5.5	4.5	2.5	1.2	-
480	8.4	7.8	7.2	6.2	4.2	2.9	1.6
490	9.6	9.0	8.4	7.4	5.4	4.1	2.9
500	13.0	12.0	11.0	10.0	8.3	7.0	5.8
510	16.0	15	14	13	11	10	9.0
520	20.0	19	18	17	15	14	13
530	21.0	20.6	19.7	19.1	16.9	15.6	14.5
540	22.6	21.7	21.4	20.3	18.3	17.1	15.8
550	24.7	24.0	23.5	22.6	20.5	19.7	18.1
560	26.8	28.2	25.6	24.5	22.5	21.3	20.0
570	27.0	26.6	26.1	25.0	23.0	21.8	20.6
580	28.1	27.7	27.2	26.1	24.1	22.6	21.5
590	26.3	25.7	25.1	24.1	22.2	20.7	18.6

under combined illumination of the sun and sky (with solar zenith angles higher than 45°) depends only on inherent optical properties of seawater. We obtained the following equation that connects spectral downward diffuse attenuation coefficient $k_{\downarrow}(\lambda)$ with the effective wavelength computed using upwelling radiation spectra in its visible part:

$$\lg k_{\downarrow}(\lambda) = A_3(\lambda) + B_3(\lambda)\lambda_{eff}, \dots 459 < \lambda_{eff} < 521 \text{ nm}, \quad (N = 49, r - \text{see Tab. 3}), \quad (8)$$

here $k_{\downarrow}(\lambda)$ is downward diffuse attenuation coefficient averaged over upper sea layer in m^{-1} . Coefficients A_3 and B_3 are given in Table 3. In all considered range of wavelengths, from 410 to 580 nm, correlation coefficients are very high with maximum value of $r = 0.98$ at $\lambda = 430 \div 440$ nm. Comparison of additional measurements, not included in derivation of Eq. (8), showed a satisfactory agreement with the prediction given by Eq. (8).

Integral downward diffuse attenuation coefficient is connected with $z_Q(1\%)$, the depth at which quantum irradiance Q , given by Eq. (2), is attenuated 100 times, by the following equation:

Table 2a. Numerical values of coefficient A_2 in Eq. (7).

$\lambda_n \setminus \lambda_k$	410	420	430	440	450	460	470
470	488.2	484.6	483.3	482.7	481.7	481.8	-
480	486.7	483.3	482.1	481.1	479.5	478.3	482.3
490	487.0	484.3	483.3	483.2	483.1	483.7	486.4
500	488.8	486.3	485.8	486.1	486.5	487.4	490.0
510	491.7	498.8	489.6	490.3	491.4	492.7	495.1
520	494.0	492.5	492.4	493.2	494.4	495.6	497.7
530	494.6	493.1	493.2	493.8	495.0	496.1	497.9
540	494.9	493.8	493.5	494.2	495.4	496.4	498.1
550	496.1	495.0	495.0	495.6	486.7	497.3	499.2
560	496.7	495.7	495.7	496.2	497.4	498.4	499.6
570	497.7	496.4	496.4	497.2	498.3	499.2	500.5
580	501.9	500.9	501.0	501.8	503.2	504.3	505.8
590	506.2	505.1	505.2	506.4	508.0	509.5	510.0

Table 2b. Numerical values of coefficient B_2 in Eq. (7).

$\lambda_n \setminus \lambda_k$	410	420	430	440	450	460	470
470	83.3	98.9	115.3	153.7	238.7	352.7	-
480	80.0	90.7	102.7	131.9	198.4	286.9	372.7
490	68.1	74.9	84.0	98.1	121.4	139.5	139.2
500	65.1	68.5	74.5	85.7	107.2	126.9	148.4
510	56.3	59.2	62.9	69.6	82.9	98.2	103.4
520	47.3	49.6	51.9	55.9	63.7	69.2	75.5
530	44.5	46.3	47.5	51.0	57.7	62.3	66.7
540	42.5	44.5	45.7	48.2	53.6	57.1	61.0
550	39.4	40.7	41.8	43.6	48.0	49.5	53.4
560	36.6	37.7	38.5	40.3	43.6	45.8	47.9
570	36.3	36.9	37.3	39.4	42.7	44.8	46.7
580	35.2	35.8	36.3	37.9	40.8	42.7	44.5
590	36.4	36.8	37.6	39.4	42.0	44.4	42.7

$$k_{\downarrow}(\lambda) = 4.61 / z_Q(1\%). \quad (9)$$

This depth, also known as the depth of euphotic zone, was proposed by Jerlov [8] in his optical classification of water masses. This value is connected with z_{90} [4] by the following relationship:

$$z_{90} = \frac{1}{4.61} z_Q(1\%). \quad (10)$$

Using experimental data taken from Refs. [11-13] we obtained the following connection between euphotic zone depth and effective wavelength:

$$z_Q\{1\% \} = 657.6 + 1.26 \lambda_{eff}, \quad (464 < \lambda_{eff} < 521 \text{ nm}), \quad (N = 80, r = 0.96). \quad (11)$$

The data set used to obtain Eq. (11) is shown in Fig. 1.

3. CONNECTION BETWEEN EFFECTIVE WAVELENGTH OF UPWELLING RADIATION AND INHERENT OPTICAL PROPERTIES OF SEAWATER

One of the important results of our measurements is a discovery of connection between effective wavelength λ_{eff} and inherent optical properties of the upper layer of the sea. Transfer and distribution of light energy in sea depth is determined by absorption and scattering processes. The accurate description of these processes is very important for overall solution of radiative transfer problems [14-16]. The in-situ experiments often lack complete set of direct measurements of inherent optical properties, especially absorption coefficient $a(\lambda)$ and phase function of

Table 3. Numerical values of coefficients A_3 , B_3 , A_4 , and B_4 in Eqs. (8) and (12).

λ, nm	$-A_3(\lambda)$	$B_3(\lambda) \cdot 10^2$	r	$-A_4(\lambda)$	$B_4(\lambda) \cdot 10^2$	r
410	13.24	2.52	0.97	13.91	2.64	0.97
420	13.23	2.52	0.97	13.81	2.62	0.98
430	13.18	2.51	0.98	13.77	2.61	0.98
440	13.06	2.48	0.98	13.48	2.55	0.98
450	12.49	2.36	0.97	12.83	2.41	0.98
460	12.39	2.38	0.97	12.57	2.35	0.97
470	12.19	2.29	0.97	12.45	2.32	0.97
480	11.77	2.20	0.97	11.87	2.20	0.97
490	11.33	2.11	0.97	11.30	2.08	0.96
500	10.43	1.93	0.96	10.41	1.90	0.96
510	8.95	1.63	0.95	8.85	1.60	0.95
520	8.11	1.46	0.94	7.83	1.39	0.94
530	7.70	1.38	0.93	7.42	1.31	0.93
540	7.07	1.25	0.92	6.74	1.17	0.92
550	6.30	1.10	0.90	5.93	1.01	0.91
560	5.60	0.96	0.88	5.20	0.86	0.89
570	5.12	0.87	0.86	4.84	0.80	0.88
580	4.62	0.78	0.84	4.28	0.70	0.85
590	-	-	-	4.30	0.72	0.86

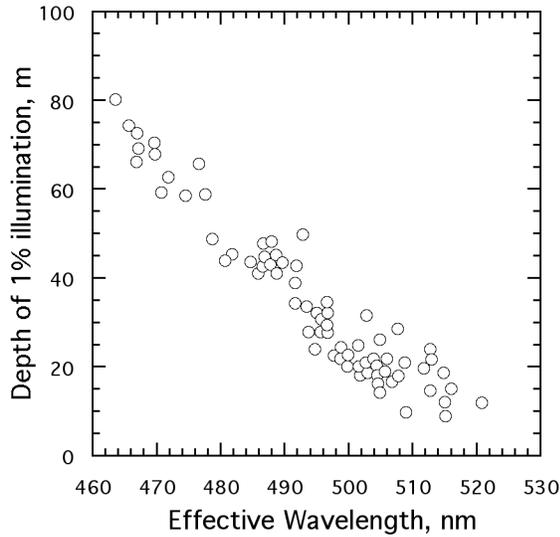


Figure 1. Dependence of euphotic zone depth (depth of 1% irradiance) on effective wavelength of upwelling radiance from the sea.

scattering $p(\lambda, \vartheta)$, where ϑ is a scattering angle. This lead to the development of indirect methods to estimate absorption coefficient and phase function of scattering [17]. The method to estimate absorption coefficient is based on a dependence of spectral signatures of upwelling light with spectral dependence of absorption coefficient.

Our analysis of experimental data showed that dependence between absorption coefficient and effective wavelength may be represented as a linear regression in semi-logarithmic scale:

$$\lg a(\lambda) = A_4(\lambda) + B_4(\lambda)\lambda_{eff}, \quad (12)$$

here $a(\lambda)$ (in m^{-1}) is an absorption coefficient averaged over subsurface layer. Coefficients A_4 and B_4 are given in Tab. 3. Spectral dependence of correlation coefficient for Eq. (12) is similar to the same of Eq. (8) and reaches maximum $r = 0.98$ at $\lambda = 420 \div 450$ nm. Majority of published data on absorption coefficient satisfy Eq. (12).

We should also note that Eqs. (8) and (12) contain values of diffuse attenuation and absorption coefficients averaged over upper sea layer without any precise information on vertical stratification of these quantities.

If we have measurements of vertical profiles of spectral beam attenuation coefficient, we can compute the following spectral depth using following definition:

$$z^*(\lambda) = \int_0^{\infty} z \exp\left[-\int_0^z c(\lambda, z') dz'\right] dz \Big/ \int_0^{\infty} \exp\left[-\int_0^z c(\lambda, z') dz'\right] dz, \quad (13)$$

here $c(\lambda, z)$ is a spectral depth profile of attenuation coefficient in m^{-1} , z is a depth in m. For a homogeneous sea $z^*(\lambda) = 1/c(\lambda)$, *i. e.* z^* coincides with spectral optical length of seawater.

The relationship between optical thickness and beam attenuation coefficient allow us to define beam attenuation coefficient averaged over upper layer of the sea using the following definition:

$$\bar{c}(\lambda) = \int_0^{\infty} \exp\left[-\int_0^z c(\lambda, z') dz'\right] dz \Big/ \int_0^{\infty} z \exp\left[-\int_0^z c(\lambda, z') dz'\right] dz. \quad (14)$$

Experimental measurements show that spectral depth $z^*(\lambda)$ has a distinctive spectral shape that makes it

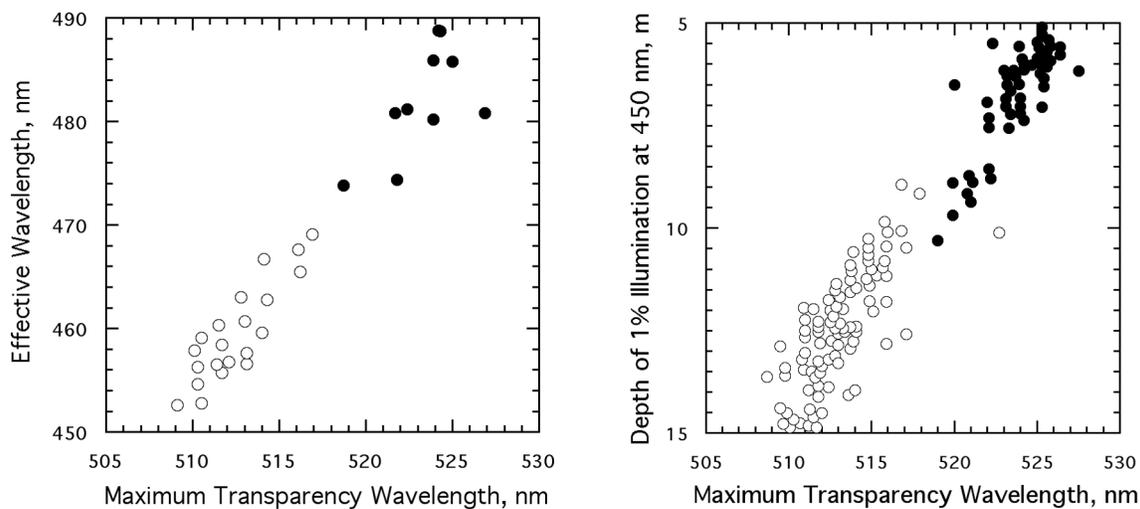


Figure 2 . Dependence of effective wavelength of upwelling sea radiance $\lambda_{eff}(\Phi)$ (left) and of the depth at which light attenuates e times $z^*(450)$ (right) with the wavelength of maximum transmission of water $\lambda_{eff}(z^*)$.

suitable to be described by the following integral parameter:

$$\lambda_{eff}(z^*) = \int_{\lambda_1}^{\lambda_2} \lambda z^*(\lambda) d\lambda / \int_{\lambda_1}^{\lambda_2} z^*(\lambda) d\lambda. \quad (15)$$

Comparison of in situ measurements made in Black and Mediterranean seas showed that $\lambda_{eff}(z^*)$ correlates both with $z^*(\lambda)$ and $\lambda_{eff}(\Phi)$, where Φ is any one of L_{\uparrow} , E_{\uparrow} , R , or r (See. Fig. 2). We do not give here regressional relationships between above mentioned parameters because a preliminary nature of these investigations.

4. CONNECTION BETWEEN EFFECTIVE WAVELENGTH OF UPWELLING RADIATION AND CHLOROPHYLL CONCENTRATION IN SEAWATER

Biological components of natural waters, the phytoplankton and products of its life cycle, play significant role in processes of absorption and scattering of radiation in visible range of spectrum. These processes cause the correlation between chlorophyll content [2] which is a dominant component in seawater and parameters of upwelling radiation. The regressional relationships similar to the following,

$$C_{chl} = A_5 J(\lambda_n, \lambda_k)^{B_5}, \quad (16)$$

are widely used [18]. here A_5 and B_5 depend on the position of wavelengths λ_n and λ_k in spectrum. When one or both these wavelength happened to be at the end of decreasing part of the spectrum, it is very difficult to obtain accurate values of color index J , and, consequently, accurate values of C_{chl} . This problem was discussed by Hovis et. al. [19]. The authors of Ref. [19] in order to increase accuracy of C_{chl} determination varied values of λ_n and λ_k using different combination of 443, 520, 550, and 670 nm channels of Coastal Zone Color Scanner.

The use of effective wavelength $\lambda_{eff}(\Phi)$ eliminates this shortcoming because this quantity is determined by the most energetic components of the optical spectrum $\Phi(\lambda)$. Our measurements showed that the effective wavelength very well correlates with the concentration of subsurface chlorophyll-*a*. The regression between C_{chl} (in mg / m^3) and effective wavelength (in nm) is expressed as follows:

$$C_{chl} = \exp\left[0.117(\lambda_{eff} - 498.2)\right], \quad (460 < \lambda_{eff} < 520 \text{ nm}), \quad (N = 74, r = 0.91). \quad (17)$$

To derive Eq. (17) we used both our data obtained in Black and Mediterranean seas and data obtained from Refs. [11, 20].

Determination of chlorophyll concentration is also possible on the basis of connection between C_{chl} and downward diffuse attenuation coefficient k_{\downarrow} . Spectral signature of downward diffuse attenuation coefficient is used by Smith and Backer [10] in their remote sensing model to derive chlorophyll concentration. Determination of k_{\downarrow} using λ_{eff} may be applied to this model and converted to simple engineering equation:

$$k_{\downarrow}(500) = \exp\left[0.0444(\lambda_{eff} - 540.4)\right]. \quad (18)$$

Also, by using Eq. (8) with coefficients taken from Tab. 3, and relationships $k_{\downarrow}(500) = 1.15/z_s$, and $\lg C_{chl} = 2.37 - 2.53 \lg(z_s)$ borrowed from Refs. [6, 21], we can obtain the following formula:

$$C_{chl} = \exp\left[0.112(\lambda_{eff} - 495.3)\right]. \quad (19)$$

The coefficients in Eq. (19) are very close to the coefficients of Eq. (17). This circumstance proves validity of dependencies (8) and (17) that connect λ_{eff} with corresponding bio-optical properties.

The level of biological productivity of waters is closely connected with a photosynthesis reaction which consists of absorption of light quantum by a phytoplankton cell to produce organic matter. As a consequence the depth profile of quantum irradiance in productive waters differs significantly from similar profile in oligotrophic waters of open ocean. This difference may be described by Lorenzen formula [22] that connects euphotic zone depth (the depth at which quantum irradiance decreases 100 times) with the subsurface concentration of chlorophyll:

$$\lg z(1\%) = 3.457 - 0.297 C_{chl}. \quad (20)$$

A combination of Eqs. (11) and (20) allows to obtain values of chlorophyll that are close (but somewhat underestimated) than values produced by Eq. (17). It is necessary to note that Eq. (20) is valid only for very turbid waters with $z_Q(1\%) \leq 30 \text{ m}$ (corresponding values of $\lambda_{eff} \geq 500 \text{ nm}$).

5. CONCLUSION

The experimentally derived relationships proposed in this paper proves usefulness of the universal integral parameter of multispectral optical signature of the ocean water, the effective wavelength λ_{eff} . By using this parameter it is possible to significantly diminish flow of primary data that characterized effectively multispectral signatures of the ocean.

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